#### Criteria for elastic fracture

When the non-linear material process zone near the crack tip is small compared to all other relevant length scales in the problem / structure, a K-annulusexists and linear elastic fracture mechanis (LEFM) is applicable. This situation is referred to as small scale yielding (SSY).

Then what are the conditions required for fracture? Discussions are focused on Mode I here while relevant concepts apply for other modes.

### Stress intensity handbooks

Under Mode I conditions, there is an equivalence between  $K_{\texttt{z}}$  and  $\zeta$ :  $G = \frac{k_{\rm L}^2}{F}$ 

One can use either  $G$  or  $K_{\mathcal{I}}$  to characterize the loading that leadsto fracture.



60

\nFor example, 
$$
\frac{1}{\sqrt{2}}
$$
 and  $\frac{1}{\sqrt{2}}$  is a specific function.

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There are also a number of ASTM standard tests such as compact tension (see Fracture Mechanics by A.T. Zehnder, Springer. 2012).

### Initiation under monotonic slow loading

Under ssy situations fracture occurs when

$$
G = G_c
$$
 or  $K_{I} = K_c$ .  
fracture toughness

The Griffith's problem is a through crack in a glass plate under tension. Assuming plane stress, G for this problem is

$$
K_{I} = \frac{1}{\sqrt{2\alpha}}
$$
  
\n $K_{I} = \frac{1}{\sqrt{2\alpha}}$   
\n $\frac{1}{\sqrt{2\alpha}}$   
\n $\frac{1}{\sqrt{2\alpha}}$   
\n $\frac{1}{\sqrt{2\alpha}}$   
\n $K_{C} = (E G_{c})^{\frac{1}{2}}$   
\n $K_{C} = (E G_{c})^{\frac{1}{2}}$ 

 $\bigcirc$ Griffith used glass taken from test tubes blown into thin-walled spheres and cylinders. A glass cutter was wed to introduced through cracks in the test sample, whoch was farther annealed to elaminate any residual stresses due to cutting.



. Temperature effect: For many materials such as metals and polymers, Kc is temperature dependent



. Thickness effect: For small thickness, there is a "loss of constraint" the plastic zone near the crack tip



ASTM condition for a valid Kte measurement

$$
t \ge 2.5 \left(\frac{K_{Tc}}{s_y}\right)^2 \le 16 r_y
$$
 ensures plane strata.

$$
Min(a, b, t.) \geq 25 \left(\frac{K_{Ic}}{G_{J}}\right)^2 \text{ ensures } S_{J}.
$$

by 300MPa

 $r_{\rm g} \sim 2.8$ mm

. Slow loading: Time to load to failure is much larger than the time for a stress wave to propagate through the material Specimen.  $\left(\frac{1}{k}\right)$  $\frac{M_{\alpha} (a, b, c)}{C}$ ,  $C = \left(\frac{E}{C}\right)$ t Speed of sound

 $\tilde{k}$ -dependency is also of importance in viscoelastic, poroelastic materials.

. Mixed mode loading: If  $G = G_c$  is valid,  $G = G_c$  represents an ellipsoidal "fracture surface" in KI, KI, KI space.  $\mathcal{G} = \frac{k_{\rm r}}{\Gamma'} + \frac{k_{\rm r}^2}{\Gamma} + \frac{k_{\rm r}^2}{2\mu} = \mathcal{G}_c$ 



Mode dependent critical fracture energy is especially important in the failure of bimaterial interfaces. Furthermore, crack loaded predominantly in Mode I tend to kink out of the craek plane. The mechanism is attributed to differences in the development of the yield/process zone in different mode (for ductile metals) or frictional effects (absent in model but present in model  $\ell$  II for brittle interfaces).

The simple criteria above let us determine whether a crack will grow or not, but they do not tell us anything about how fact, how far, or in what direction the crack will grow. These topics are to be address shortly.

## Crack growth stability (how fan?)

We have shown in the analysis of the DCB speciment that  $G\propto a^2$  for fixed lad and  $\zeta$  a  $a^{-4}$  for fixed displacement. The crack growth is likely stable in the latter case since  $\frac{dG}{da} < 0$ . There are additional stablizing mechanisms.

Loading by compliant systems



 $\bigcirc$ 

 $p = cq = cs q_s$ ,  $q_t = q + q_s \rightarrow q_t = (\frac{1}{c} + \frac{1}{c_s}) \rho$ Note that  $C_S \rightarrow \infty$   $\Leftrightarrow$  fixed displacement

$$
\begin{aligned}\n\left| \int_{\mathfrak{g}_{k}} \frac{1}{2} C \varphi^{2} + \frac{1}{2} C_{5} \varphi_{s}^{2} &= \frac{1}{2} \frac{\rho^{2}}{C} + \frac{1}{2} \frac{\rho^{2}}{C_{5}} = \frac{1}{2} \left( \frac{1}{C} + \frac{1}{C_{5}} \right) \frac{\varphi_{s}^{2}}{C_{5}} / \left( \frac{1}{C} + \frac{1}{C_{5}} \right)^{2} &= \frac{1}{2} \frac{\varphi_{s}^{2}}{C_{5}} / \left( \frac{1}{C} + \frac{1}{C_{5}} \right)\n\end{aligned}
$$
\n
$$
\left| \int_{\mathfrak{g}_{k}} \frac{1}{2} \frac{1}{
$$

You may find this system is also eventually stable as  $l$  increases ( $\rho\psi$ ).

Resistance curve (R-curve)

In some cases, crack growth can be stable even when  $\frac{\partial G}{\partial k} > 0$ . How so? For brittle materials.  $G_c$  is a constant. For ductile materials (particular thin sheets). It is generally found that  $G_c$  increases as the crack grows.



Mixed-mode fracture initiation & growth (Which direction?)



Subject to a combination of Mode I and Mode II headings crack will generally not propagate straight ahead (unless  $\Theta = 0$  is a weak interface).

Our previews analysis has been on  $9=-\frac{\partial w}{\partial a}|_{\theta=0}$ . What if  $\theta=0$ ? There are a few theories such as  $\sigma$  maximum hoop strees  $\varnothing$  maximum energy release rate.

(62) O Maximum hoop stress theory: Grack will grow in the direction  $\partial_{\mathbf{x}}$  of maximum hoop stress, when  $\sqrt{\Gamma}$   $\zeta_{\theta\theta}$  (r,  $\theta_*$ )  $\geq$  Constant.

$$
\rightarrow \frac{\partial \phi_{\theta\theta}}{\partial \theta} \bigg|_{\theta_{\mathcal{H}}} = 0, \quad \frac{\partial^2 \phi_{\theta\theta}}{\partial \theta^2} \bigg|_{\theta_{\mathcal{H}}} < 0 \quad , \quad \phi_{\theta\theta}(r, \theta_{\mathcal{H}}) \geq \frac{K_{TC}}{\sqrt{2\pi r}}
$$

For this problem 
$$
K_{I} = 6\sqrt{\pi a} \sin \beta
$$
,  $K_{I} = 6\sqrt{\pi a} \cos \beta$ 

$$
\measuredangle_{\theta\theta} (r,\theta) = \frac{1}{\sqrt{2\pi r}} \left( \frac{1}{K_{\perp}} \cos{\frac{\theta}{2}} \frac{r + \cos{\theta}}{2} - \frac{1}{K_{\perp}} \frac{3}{2} \sin{\theta} \cos{\frac{\theta}{2}} \right) + O(r^{\circ})
$$



Under pure Mode  $I(\xi = o^*)$ , the crack will grow at  $\theta_k = -70.6$  at a stress intensity factor of  $K_{I\!I}$  = 0.87  $K_{I\!C}$ .  $I$  Allowing Kinking leads to a slightly smaller  $K_{\rm I\!I} \int$ 

2 Maximum energy release rate criterion:  $\frac{\partial 6}{\partial \theta}|_{\theta_{\mathbf{x}}} = 0$ ,  $\frac{2^2 6}{\partial \theta^2}|_{\theta_{\mathbf{x}}} < 0$ ,  $\int (\theta_{\mathbf{x}}) = 6$  $\Rightarrow$   $\theta_x = -75.6^\circ$  at  $K_{\mathbb{I}} = 0.817 K_{\mathbb{I}}c$  under pure Mode I loading.



 $(63)$ 

Why  $-70^\circ$  not  $-45^\circ$ ? The presence of crack disturbs the stress field. The crack path will evolve to lie along the plane of maximum principle stress.

### Mode I cyclic loading Fatigue

Under cyclic loading, materials can fail due to fatigne at stress levels well below their static strength.



# Experimental methods

In some cases, it may not be pratical to determine stress intensity factors from analytical or computational methods. For example, perhaps the loading is unknown, or is dynamic. One may wish to determine the SIF experimentally, based on local measurement of stress, strain and displacement. There are a number of experimental methods that have been developed, including photoelasticity, interferometry, digital image correlation, strain guages and so on (see Chapter 5 in Fracture Mechanics by A.T. Zehnder)

 $\bigcirc$ 

Only strain guages method by Dally and Stanford (1987) is discussed here.



 $K_{\tau}$ , Ao, A $v_2$  – Three parameters, but only one measurement. What do you do?